Integral Equations for Wave Scattering: Numerical Solution and Wavenumber-Explicit Bounds

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Joint work with collaborators, notably: Ivan Graham, Euan Spence (Bath), Steve Langdon (Reading)

Recent Advances in the Numerical Analysis of PDEs: Celebrating the 65th birthday of Prof Ivan Graham This talk is about what, with assistance of Ivan, I've worked on throughout my career, namely $% \left({{{\mathbf{r}}_{\mathrm{s}}}_{\mathrm{s}}} \right)$

(i) solving

$$\Delta u + \frac{k^2 u}{k} = 0$$

by integral equation methods.

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And what I've worked on for much of my career, partly collaborating with Ivan

(ii) understanding how everything depends on k.

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I'll talk today about two specific acoustics problems.

Problem 1



Problem 2



Problem 1: Sound Propagation over Inhomogeneous Impedance Plane

Problem I've worked on since 1985, 11 papers in *J Sound Vib*, *IMA J Numer Anal*, *Math Meth Appl Sci*, *J Math Anal Appl*, *Numer Math*, *Proc R Soc A*, *SINUM*, the last two with **Steve Langdon** (Ivan's 1999 PhD student).



- k > 0 is the wavenumber
- $\beta \in L^\infty(\mathbb{R})$, the **impedance**, is typically piecewise constant, and $\Re \beta \ge 0$
- u satisfies standard radiation condition at infinity

Problem 1: Simplest Case, $\beta \equiv 0$



Solution by method of images is

$$u(x) = G_0(x, y) := \frac{i}{4} H_0^{(1)}(\mathbf{k}\mathcal{R}) + \frac{i}{4} H_0^{(1)}(\mathbf{k}\mathcal{R}')$$

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where $H_0^{(1)}$ is a Hankel function. N.B. $G_0(x, y)$ is just the **Neumann Green's function**.



By Green's theorem – $G_0(\boldsymbol{x},\boldsymbol{y})$ is the Neumann Green's function –

$$u(x) = G_0(x, y) + \int_{\gamma} G_0(x, z) \mathbf{i} \mathbf{k} \beta(z) u(z) \mathrm{d} s(z), \quad x \in \overline{U}.$$



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In particular, where $\phi(x_1) := u((x_1, 0))$, and if $y = (0, y_2)$,

$$\phi(x_1) = \frac{i}{2} H_0^{(1)} \left(\frac{k}{\sqrt{x_1^2 + y_2^2}} \right) - \frac{k}{2} \int_a^b H_0^{(1)}(\frac{k}{x_1 - z_1}) \beta(z_1) \phi(z_1) dz_1, \quad a \le x_1 \le b.$$

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$$\phi(t)=f(t)+\int_0^1\kappa(t-s)\phi(s)\mathrm{d} s,\quad 0\leq t\leq 1,$$

and $f \in L_1[0,1]$ and κ is in a Nikol'skii space intermediate between $L_1[-1,1]$ and $W_1^1[-1,1]$, then, for $m \in \mathbb{N}$,

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with $s(x) = \int_0^x \kappa(t) dt = O(x \ln |x|)$ and $\psi \in W_1^1[0, 1]$. See C-W, Gover (1989).

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Note that in the case of equations (1) and (2) mesh grading is hardly going to make any difference at all to the convergence rates (using piecewise (anstants) so why not use (uniform meshes and use a fast method for solving the linear systems? Good hick with your work Rease Keep me a your wailing list! Yvan /Ivan Graham.

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In C-W, Rahman, Ross (2002) we followed this advice to the letter!

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Numerical method ingredients:

- Uniform grid, stepsize h = (b a)/N
- Approximate $\beta \in L^{\infty}[a,b]$ by its local average on each grid subinterval
- Piecewise constant approximation for ϕ
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Theorem (C-W, Rahman, Ross 2002: plane wave incidence)

Suppose β takes values in compact subset Q of the right-hand complex plane. Then, after 7 iterations, provided $kh \leq c_Q$,

 $\|\phi - \phi_h\|_{\infty} \le C_Q k h \log(1/kh),$

and ϕ_h is computed in $O(N \log N)$ operations.

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This my first k-explicit estimate, and an early k-explicit estimate for BEM for wave scattering (cf. Löhndorf, Melenk 2011, Graham, Löhndorf, Melenk, Spence 2015).

Problem 2: Scattering by Sound Soft Obstacles: Integral Equations and k-Explicit Bounds on the Operators



Assume throughout that Ω_{-} is bounded and Lipschitz.

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$$\mathcal{U}_{+} u^{\text{inc}} \qquad \Delta u + k^{2}u = 0$$

$$\Omega_{-} \qquad \Gamma \ u = 0$$

$$u - u^{\text{inc}} \text{ satisfies radiation condition}$$

$$\Omega_{+}$$

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Theorem (Green's Representation Theorem)

$$u(x) = u^{\rm inc}(x) + \int_{\Gamma} \Phi(x,y) \partial_n^+ u(y) \, ds(y), \quad x \in \Omega_+.$$

where

$$\Phi(x,y) := \frac{\mathrm{i}}{4} H_0^{(1)}(\mathbf{k}|x-y|) \quad \text{(2D),} \quad := \frac{1}{4\pi} \frac{\mathrm{e}^{\mathrm{i}\mathbf{k}|x-y|}}{|x-y|} \quad \text{(3D).}$$



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Taking a linear combination of Dirichlet (γ_+) and Neumann (∂_n^+) traces, we get the **boundary integral equation** (Burton & Miller 1971)

$$\frac{1}{2}\partial_n^+ u(x) + \int_{\Gamma} \left(\frac{\partial \Phi(x,y)}{\partial n(x)} + \mathrm{i}\eta \Phi(x,y) \right) \partial_n^+ u(y) ds(y) = f(x), \quad x \in \Gamma,$$

where

$$f := \partial_n^+ u^{\rm inc} + {\rm i}\eta\gamma_+ u^{\rm inc}.$$

$$\begin{split} & \underbrace{\lambda u}_{u^{\mathrm{inc}}} & \Delta u + k^2 u = 0 \\ & & \Gamma \ u = 0 \\ & & u^{\mathrm{inc}} & u - u^{\mathrm{inc}} \text{ satisfies radiation condition} \\ & & \frac{1}{2} \partial_n^+ u(x) + \int_{\Gamma} \left(\frac{\partial \Phi(x,y)}{\partial n(x)} + \mathrm{i} \eta \Phi(x,y) \right) \partial_n^+ u(y) ds(y) = f(x), \quad x \in \Gamma, \\ & \text{in operator form } A_{k,\eta} \partial_n^+ u = f := \partial_n^+ u^{\mathrm{inc}} + \mathrm{i} \eta \gamma_+ u^{\mathrm{inc}}. \end{split}$$

$$\begin{split} & \mathcal{M}_{\mathbf{x}} \ u^{\mathrm{inc}} \qquad \Delta u + \mathbf{k}^2 u = 0 \\ & \mathbf{\Gamma} \ u = 0 \\ & \mathbf{\Omega}_{-} \qquad u - u^{\mathrm{inc}} \text{ satisfies radiation condition} \\ & \frac{1}{2} \partial_n^+ u(x) + \int_{\Gamma} \left(\frac{\partial \Phi(x, y)}{\partial n(x)} + \mathrm{i} \eta \Phi(x, y) \right) \partial_n^+ u(y) ds(y) = f(x), \quad x \in \Gamma, \\ & \mathrm{in \ operator \ form \ } A_{\mathbf{k}, \eta} \partial_n^+ u = f := \partial_n^+ u^{\mathrm{inc}} + \mathrm{i} \eta \gamma_+ u^{\mathrm{inc}}. \end{split}$$

Theorem (Burton & Miller 1971, Mitrea 1996, C-W & Langdon 2007)

If $\eta \in \mathbb{R}$, $\eta \neq 0$, then this integral equation is uniquely solvable in $L^2(\Gamma)$.

$$\begin{split} & \mathcal{L}_{u} u^{\text{inc}} \qquad \Delta u + k^{2}u = 0 \\ & \mathbf{\Gamma} \quad u = 0 \\ & \mathbf{u} - u^{\text{inc}} \text{ satisfies radiation condition} \\ & \frac{1}{2}\partial_{n}^{+}u(x) + \int_{\Gamma} \left(\frac{\partial \Phi(x,y)}{\partial n(x)} + i\eta \Phi(x,y)\right) \partial_{n}^{+}u(y)ds(y) = f(x), \quad x \in \Gamma, \\ & \text{erator form } A_{u} \partial^{+}u = f := \partial^{+}u^{\text{inc}} + in\gamma_{u}u^{\text{inc}} \end{split}$$

in operator form $A_{\mathbf{k},\eta}\partial_n^+ u = f := \partial_n^+ u^{\mathrm{inc}} + \mathrm{i}\eta\gamma_+ u^{\mathrm{inc}}$.

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Theorem (C-W & Monk 2008, C-W, Graham, Langdon, Lindner 2009, Han, Tacy, Galkowski 2015, Baskin, Spence, Wunsch 2016)

If $\eta \approx \mathbf{k}$ and Ω_{-} is: (i) star-shaped with respect to a ball and piecewise smooth; or (ii) C^{∞} and non-trapping; then, as an operator on $L^{2}(\Gamma)$, for $\mathbf{k} \geq k_{0}$,

 $\|A_{\boldsymbol{k},\eta}^{-1}\| \lesssim 1, \quad \|A_{\boldsymbol{k},\eta}\| \lesssim \boldsymbol{k}^{1/2}\log \boldsymbol{k}, \quad \operatorname{cond} A_{\boldsymbol{k},\eta} \lesssim \boldsymbol{k}^{1/2}\log \boldsymbol{k}.$

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| k | $\ A_{k,k}\ $ | p | $\ A_{k,k}^{-1}\ $ |
|-----|---------------|------|--------------------|
| 5 | 3.089 | | 1.024 |
| 10 | 3.611 | 0.23 | 1.024 |
| 20 | 4.608 | 0.35 | 1.023 |
| 40 | 6.032 | 0.39 | 1.023 |
| 80 | 8.117 | 0.43 | 1.023 |
| 160 | 11.068 | 0.45 | 1.023 |
| 320 | 15.253 | 0.46 | 1.023 |
| 640 | 21.177 | 0.47 | 1.023 |

Galerkin BEM discretisations for a square of sidelength 2, using piecewise constants, 10 elements per wavelength, from Betcke, C-W, Graham, Langdon, Lindner (2011)





Step 1 (Resolvent Estimate). Show that, for every R > 0, if g = 0,

 $\|\nabla u\|_{L^{2}(\Omega_{R})} + \mathbf{k}\|u\|_{L^{2}(\Omega_{R})} \lesssim c(\mathbf{k})\|f\|_{L^{2}(\Omega_{+})},$

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Step 2 (DtN Map Bound). It follows that, if f = 0,

 $\|\partial_n^+ u\|_{L^2(\Gamma)} \lesssim c(\mathbf{k}) \left(\|\nabla_{\Gamma} g\|_{L^2(\Gamma)} + \mathbf{k} \|g\|_{L^2(\Gamma)} \right)$



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Step 3 As (C-W, Graham et al 2012) $A_{k,k}^{-1} = I - (P_{DtN}^+ - ik)P_{ItD}^-$ and P_{ItD}^- is bounded in Spence (2015), it follows that $||A_{k,k}^{-1}|| \lesssim c(k)k^{1/2}$.



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Sharper by a factor $k^{1/2}$ if Ω_{-} is starlike (Melenk 1995), or is C^{∞} (Baskin et al 2016)

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$$\Omega_{-} \xleftarrow{a} \Omega_{-}$$

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Theorem (C-W, Graham, Langdon, Lindner 2009, C-W, Spence 2017)

For trapping domains like this (neutrally trapping) it holds, for $k \ge k_0$, that

$$\|A_{m k,m k}^{-1}\|\lesssim m k^{5/2},$$
 and that $m k^{9/10}\lesssim \|A_{m k,m k}^{-1}\|$

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$$\|A_{{m k},{m k}}^{-1}\| \lesssim {m k}^{5/2},$$
 and that ${m k}^{9/10} \lesssim \|A_{{m k},{m k}}^{-1}\|$

if ka is a multiple of π .

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if ka is a multiple of π .

| k | $\ A_{k,k}\ $ | p | $ A_{k,k}^{-1} $ | p |
|-----|---------------|------|--------------------|------|
| 5 | 4.835 | | 1.969 | |
| 10 | 5.201 | 0.11 | 3.121 | 0.66 |
| 20 | 5.629 | 0.11 | 5.539 | 0.83 |
| 40 | 6.182 | 0.14 | 10.322 | 0.90 |
| 80 | 8.112 | 0.39 | 19.774 | 0.94 |
| 160 | 11.066 | 0.45 | 38.351 | 0.96 |
| 320 | 15.254 | 0.46 | 75.156 | 0.97 |

Galerkin BEM discretisations for the trapping domain with ka a multiple of π , using piecewise constants, 10 elements per wavelength, from Betcke, C-W, Graham, Langdon, Lindner (2011)

Conclusions

I've said something about integral equation formulations and k-explicit estimates for two (representative) problems in acoustics that have connected me to Ivan over the last 31 years.

Problem 1



Conclusions

Ivan thank you for your friendship over 31 years ... and Happy 65th Birthday!





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